

INTRO TO MODERN COSMOLOGY, 2nd ED.
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PROBLEM SOLUTIONS BY BEN DANIEL

CHAPTER 2

2.1 $M_{\text{GAL}} \approx 10^{11} M_{\odot} \approx 10^{11} \times 2 \times 10^{30} \text{ kg} = 2 \times 10^{41} \text{ kg}$

$V_{\text{GAL}} \approx (1 \text{ Mpc})^3 \approx (3 \times 10^{16} \times 10^6 \text{ m})^3 = 3 \times 10^{47} \text{ m}^3$

$\rho_{\text{UNIV}} = M_{\text{GAL}} / V_{\text{GAL}} = 2 \times 10^{41} \text{ kg} / 3 \times 10^{47} \text{ m}^3 = \boxed{7 \times 10^{-27} \text{ kg/m}^3}$

$M_{\text{EARTH}} \approx 6 \times 10^{24} \text{ kg}$

$V_{\text{EARTH}} \approx \frac{4}{3} \pi (6.4 \times 10^6 \text{ m})^3 = 1.7 \times 10^{20} \text{ m}^3$

$\rho_{\text{EARTH}} = M_{\text{EARTH}} / V_{\text{EARTH}} = 6 \times 10^{24} \text{ kg} / 1.7 \times 10^{20} \text{ m}^3 = \boxed{3.5 \times 10^4 \text{ kg/m}^3}$

$\rho_{\text{EARTH}} / \rho_{\text{UNIV}} = 5 \times 10^{30}$

2.2 BECAUSE THE PECULIAR VELOCITIES ARE RANDOMLY ORIENTED, $v^2 = v_x^2 + v_y^2 + v_z^2$ AND $v_x = v_y = v_z$.

THUS THE SQUARE OF THE RADIAL PECULIAR VELOCITY IS $1/3$ OF THE SQUARE OF THE RANDOMLY ORIENTED PECULIAR VELOCITY.

$v_{\text{RAD}}^{\text{PEC}} = \sqrt{\frac{600^2}{3}} \text{ km/s} = 350 \text{ km/s}$

HUBBLE'S LAW STATES THAT $H_0 = v_{\text{RAD}} / r$. IF WE CAN TOLERATE AN ERROR OF 10% IN H_0 AND WE ASSUME THAT ALL UNCERTAINTY IN v_{RAD} IS DUE TO THE PECULIAR

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VELOCITY, v_{RAD} , THEN $v \geq \frac{v_{\text{RAD}}}{0.10} / H_0$.

a) $H_0 = 100 \text{ (km/s) / Mpc}$; $v \geq \frac{(350)}{0.10} / 100 = \boxed{35 \text{ Mpc}}$

b) $H_0 = 50 \text{ (km/s) / Mpc}$; $v \geq \frac{(350)}{0.10} / 50 = \boxed{70 \text{ Mpc}}$

2.3 IF THERE WERE A NET CHANGE, ITS MOTION WOULD GENERATE AN OVERALL MAGNETIC FIELD THAT WE DO NOT OBSERVE.

2.4 $f = E/h = \frac{(13.6 \text{ eV})(1.6 \times 10^{-19} \text{ J/eV})}{6.6 \times 10^{-34} \text{ J}\cdot\text{s}} = \boxed{3.3 \times 10^{15} \text{ Hz}}$

$T \cong E/3k_B = \frac{(13.6 \text{ eV})(1.6 \times 10^{-19} \text{ J/eV})}{(3)(1.4 \times 10^{-23} \text{ J/K})} = \boxed{5.2 \times 10^5 \text{ K}}$

2.5 $f_{\text{PEAK}} \cong 2.8 k_B T/h$. SINCE k_B AND h ARE CONSTANTS, SO IS $f_{\text{PEAK}}/T \cong 2.8 k_B/h$.

$\frac{f_{\text{PEAK}}}{T} \cong \frac{(2.8)(1.38 \times 10^{-23} \text{ J/K})}{6.6 \times 10^{-34} \text{ J}\cdot\text{s}} = \boxed{5.8 \times 10^{10} \text{ Hz/K}}$

$f_{\text{PEAK} \odot} \cong (5.8 \times 10^{10} \text{ Hz/K})(5800 \text{ K}) = \boxed{3.4 \times 10^{14} \text{ Hz}}$

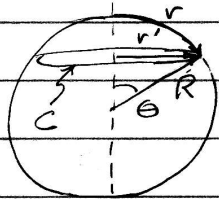
2.6 $f_{\text{PEAK, CMB}} \cong (5.8 \times 10^{10} \text{ Hz/K})(2.725 \text{ K}) = 1.6 \times 10^{11} \text{ Hz}$

$\lambda_{\text{PEAK, CMB}} = c/f_{\text{PEAK, CMB}} \cong \frac{3 \times 10^8 \text{ m/s}}{1.6 \times 10^{11} \text{ s}^{-1}} = \boxed{1.9 \times 10^{-3} \text{ m}}$

BY (2.10) AND (2.11): $\epsilon_{\text{RAD}} = \alpha T^4 = (7.565 \times 10^{-16} \text{ J/m}^3 \text{K}^4)(2.725 \text{ K})^4 = \boxed{4.17 \times 10^{-14} \text{ J/m}^3}$

CHAPTER 4

4.1



$$\sin \theta = r/R$$

RADIAN MEASUREMENT:

$$\frac{r}{2\pi R} = \frac{\theta}{2\pi} \Rightarrow \theta = r/R$$

$$C = 2\pi r' = 2\pi R \sin \theta = \boxed{2\pi R \sin(r/R)}$$

CHECKS:

$$\text{At } \theta = 0; C = 2\pi R \sin(0) = 0 \quad \checkmark$$

$$\text{At } \theta = 90^\circ; C = 2\pi R \sin(\pi/2) = 2\pi R \quad \checkmark$$

For $r \ll R$, $\sin(r/R) \rightarrow r/R$, so

$$C = 2\pi R (r/R) = \boxed{2\pi r, \text{ AS EXPECTED IN A FLAT GEOMETRY.}}$$

4.2

THE AREA OF A CIRCLE IN THE SPHERICAL GEOMETRY

ABOVE IS

$$A_s = \int_0^r C dr = 2\pi R \int_0^r \sin\left(\frac{r''}{R}\right) dr'' \quad (\text{let } u = \frac{r''}{R}; du = \frac{dr''}{R})$$

$$= 2\pi R^2 \left(-\cos\left(\frac{r''}{R}\right) \right) \Big|_0^r$$

$$= 2\pi R^2 \left(1 - \cos\left(\frac{r}{R}\right) \right)$$

$$N_s = n A_s = \boxed{2\pi n R^2 \left(1 - \cos\left(\frac{r}{R}\right) \right)}, \text{ TAYLOR EXPANDING,}$$

$$= 2\pi n R^2 \left(1 - \left(1 - \frac{r^2}{2R^2} + \dots \right) \right)$$

$$N_F \cong (2\pi n R^2) \left(\frac{r^2}{2R^2} \right) = \boxed{\pi n r^2}$$

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4.2
cont.

So, $N_S = 2\pi n R^2 (1 - \cos(\frac{r}{R}))$ AND $N_F = 4\pi r^2$

$$\frac{N_S}{N_F} = \frac{2R^2(1 - \cos(\frac{r}{R}))}{r^2} = \frac{2(1 - \cos \theta)}{\theta^2}$$

TO SHOW THE RELATIONSHIP BETWEEN N_S AND N_F REDUCES TO SHOWING THE RELATIONSHIP BETWEEN $2(1 - \cos \theta)$ AND θ^2 .

OBVIOUSLY THEY ARE BOTH 0 AT $\theta = 0$. TO INVESTIGATE THEIR BEHAVIOR FOR $\theta > 0$, TAKE THEIR DERIVATIVES:

$$\frac{d}{d\theta} (2(1 - \cos \theta)) = 2 \sin \theta \quad \frac{d}{d\theta} \theta^2 = 2\theta$$

SINCE $\sin \theta < \theta$ FOR ALL $\theta > 0$, $N_S < N_F$ FOR ALL $r > 0$, AND WE SEE FEWER GALAXIES IN A SPHERICAL UNIVERSE.

4.3

FOR A GIVEN UNIVERSE, K IS FIXED, SO $K > 0$ CAN NEVER BECOME $K < 0$ AND VICE VERSA, SO THE OPEN/CLOSED STATUS OF A UNIVERSE CAN NEVER CHANGE.

CHAPTER 5

5.1 THE UNIVERSE CAN NEVER LOSE/GAIN ENERGY
AS THERE IS NOWHERE ELSE FOR IT TO GO/COME FROM.

5.2 (2.4) $E^2 = m^2 c^4 + p^2 c^2$, BUT FOR PHOTONS, $m=0$,
SO $E^2 = p^2 c^2 \rightarrow E = pc$ OR $p = E/c$
FOR PHOTON $v=c$, SO $\vec{v} \cdot \vec{p} = (c)(E/c) = E$

SO,

$$P = \frac{1}{3} n \langle \vec{v} \cdot \vec{p} \rangle = \boxed{\frac{1}{3} n \langle E \rangle}$$

5.3 AS ALWAYS, BEGIN WITH THE FLUID EQUATION

$$\dot{\rho} + 3 \frac{\dot{a}}{a} \left(\rho + \frac{P}{c^2} \right) = 0$$

SUBSTITUTING: $P = (\gamma - 1) \rho c^2$ TO GET

$$\dot{\rho} + 3 \frac{\dot{a}}{a} (\gamma \rho) = 0$$

THIS IS THE SAME AS THE MATTER-ONLY CASE
WITH $3 \rightarrow 3\gamma$, SO $\boxed{P(a) = P_0 / a^{3\gamma}}$

THE FRIEDMANN EQUATION WITH $K=0$ THEN IS

$$\left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \rho = \frac{8\pi G}{3} \left(\frac{P_0}{a^{3\gamma}} \right)$$

THUS

$$\dot{a}^2 = \frac{8\pi G}{3} \cdot P_0 a^{2-3\gamma}$$

ASSUMING A POWER LAW SOLUTION, $a \propto t^q \Rightarrow \dot{a} \propto t^{q-1}$ AND

EQUATING POWERS GIVES $2(q-1) = q(2-3\gamma)$, OR $q = \frac{2}{3\gamma}$.

(FOR MATTER, $\gamma=1 \Rightarrow q=\frac{2}{3}$; FOR RADIATION, $\gamma=\frac{4}{3} \Rightarrow q=\frac{1}{2}$.)

⑥

5.3
Cont.

Thus $a(t) = \left(\frac{t}{t_0}\right)^{2/3\gamma}$; $p(t) = \frac{p_0}{a^{3\gamma}} = \frac{p_0 t_0^2}{t^2}$

IF $p = -p_0$ THEN $\gamma = 0$, q BLOWS UP AND THIS APPROACH GIVES NO SOLUTION. INSTEAD, RETURN TO $p = p_0/a^{3\gamma}$. HAVING $\gamma = 0$ MEANS $p = p_0$ AND $\dot{a}^2 = \frac{8\pi G}{3} p_0 a^2$. HENCE,

$\dot{a} = \pm \sqrt{\frac{8\pi G p_0}{3}} a$, A FIRST ORDER O.D.E.

WITH SOLUTION $a(t) \propto \exp\left(\pm \sqrt{\frac{8\pi G p_0}{3}} t\right)$.

5.4

$a(t) = \left(t/t_0\right)^{2/3\gamma}$; $p \propto \left(t/t_0\right)^{-2}$

SINCE K IS A CONSTANT,

K/a^2 HAS THE SAME TIME DEPENDENCE AS a^{-2} ,

THAT IS, $-4/3\gamma$. THUS, IF p AND K/a^2 HAVE THE SAME TIME DEPENDENCE,

$-4/3\gamma = -2 \Rightarrow \boxed{\gamma = 2/3}$

THE FULL FRIEDMANN EQN,

$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} p - \frac{K}{a^2}$

WITH $\gamma = 2/3$ (AND HENCE $p = p_0/a^{3\gamma} = p_0/a^2$)

$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \frac{p_0}{a^2} - \frac{K}{a^2} \Rightarrow$

$\dot{a}^2 = \frac{8\pi G p_0}{3} - K.$

SO, WITH $K < 0$, $\dot{a}^2 > 0 \Rightarrow \dot{a}$ IS A REAL

CONSTANT AND $\boxed{a \propto t}$

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S.S WITH $K > 0$; $P = 0$; $\rho = \rho_0/a^3$, TAKE $\beta = \frac{4\pi G \rho_0}{3}$.

THUS:

$$a(\theta) = \frac{\beta}{K}(1 - \cos\theta) \text{ AND } t(\theta) = \frac{\beta}{K^{3/2}}(\theta - \sin\theta)$$

$$\dot{a} = \frac{da}{dt} = \frac{da}{d\theta} \left(\frac{dt}{d\theta} \right)^{-1}$$

$$\frac{da}{d\theta} = \frac{\beta}{K} \sin\theta ; \quad \frac{dt}{d\theta} = \frac{\beta}{K^{3/2}} (1 - \cos\theta)$$

$$\frac{da}{dt} = \left(\frac{\beta}{K} \sin\theta \right) \left(\frac{K^{3/2}}{\beta(1 - \cos\theta)} \right) = \frac{\sqrt{K} \sin\theta}{(1 - \cos\theta)}$$

$$\left(\frac{\dot{a}}{a} \right)^2 = \frac{K \sin^2\theta}{(1 - \cos\theta)^2} \cdot \frac{K^2}{\beta^2 (1 - \cos\theta)^2} = \frac{K^3 \sin^2\theta}{\beta^2 (1 - \cos\theta)^4}$$

$$\left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi G \rho}{3} - \frac{K}{a^2} = \frac{8\pi G \rho_0}{3 a^3} - \frac{K}{a^2} = \frac{2\beta - Ka}{a^3}$$

$$= \frac{2\beta - K\left(\frac{\beta}{K}(1 - \cos\theta)\right)}{\left(\frac{\beta}{K}(1 - \cos\theta)\right)^3} = \frac{K^3 \beta (1 + \cos\theta)}{\beta^3 (1 - \cos\theta)^3} \cdot \frac{(1 - \cos\theta)}{(1 - \cos\theta)}$$

$$= \frac{K^3 \sin^2\theta}{\beta^2 (1 - \cos\theta)^4} \quad \checkmark$$

THIS UNIVERSE OSCILLATES WITH A PERIOD OF

$$\tau = \frac{8\pi^2 G \rho_0}{3 K^{3/2}} \quad \leftarrow \quad (\alpha(\theta) = 0 \text{ FOR } \theta = 2n\pi)$$

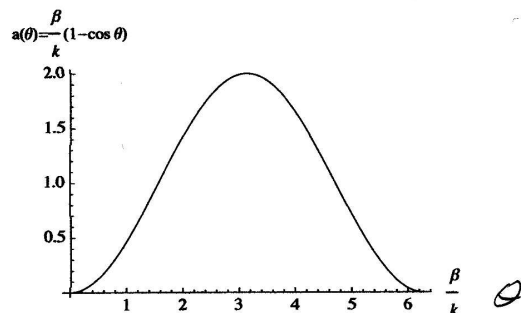
$$a_{\text{MAX}} = \frac{8\pi G \rho_0}{3 K} \quad \leftarrow$$

$$t(2n\pi) = 2n\pi \beta / K^{3/2}$$

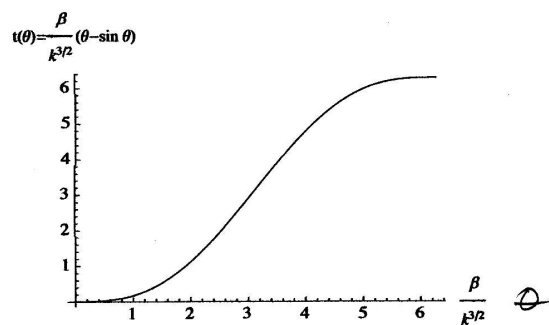
$$(\alpha_{\text{MAX}} \text{ OCCURS AT } \theta = \pi, \text{ SO } \alpha(\pi) = \alpha_{\text{MAX}} = 2\beta/K.)$$

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Plot $\left[1 - \text{Cos}[\theta], \{\theta, 0, 2 \text{ Pi}\}, \text{AxesLabel} \rightarrow \left\{\frac{\beta}{k}, "a(\theta) = \frac{\beta}{k} (1 - \text{cos } \theta)"\right\}\right]$

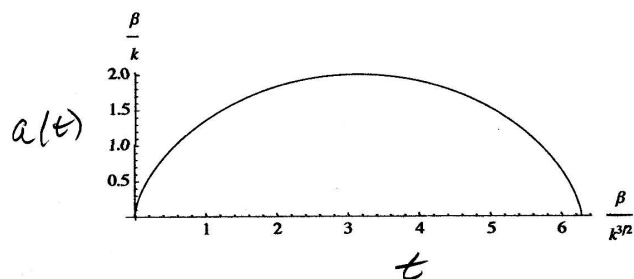


Plot $\left[\theta - \text{Sin}[\theta], \{\theta, 0, 2 \text{ Pi}\}, \text{AxesLabel} \rightarrow \left\{\frac{\beta}{k^{3/2}}, "t(\theta) = \frac{\beta}{k^{3/2}} (\theta - \text{sin } \theta)"\right\}\right]$



Below is a plot of $a(t)$ parameterized by θ . As in the plot above of $a(\theta)$, the scale factor reaches a maximum of $a_{\max} = \frac{2\beta}{k}$. This occurs at $t = \frac{\pi\beta}{k^{3/2}}$.

ParametricPlot $\left[\{\theta - \text{Sin}[\theta], 1 - \text{Cos}[\theta]\}, \{\theta, 0, 2 \text{ Pi}\}, \text{AxesLabel} \rightarrow \left\{\frac{\beta}{k^{3/2}}, \frac{\beta}{k}\right\}\right]$



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5.4 $K < 0; P = 0; \rho = \frac{\rho_0}{a^3}$

IF LAST TERM OF FRIDOMAN EQN DOMINATES,

$$\dot{a} = \sqrt{|K|} \Rightarrow a = \sqrt{|K|} t + a_0 = \sqrt{|K|} t + 1$$

$$a(t) = \sqrt{|K|} t + 1, \quad \boxed{a \propto t}$$

$$\rho(t) = \rho_0 / (\sqrt{|K|} t + 1)^3 \quad \text{so} \quad \boxed{\rho(t) \propto t^{-3}}$$

THIS IS A THINNING UNIVERSE THAT NEVER COLLAPSES. IT IS STABLE.